

D-Module Techniques for Solving Differential Equations in the Context of Feynman Integrals

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Motivation

Aim: exploit algebraic geometry behind Feynman integrals

- extraction of properties of Feynman integrals from their PDEs
- ⋄ algorithmic computation of series solutions of PDEs by algebraic methods
- evaluation of Feynman integrals
- providing a dictionary between algebraic analysis and high energy physics

Linear PDEs through an algebraic lens

Definition

The Weyl algebra is obtained from the free algebra over $\mathbb C$

$$D := \mathbb{C}[x_1,\ldots,x_n]\langle \partial_1,\ldots,\partial_n \rangle$$

by imposing the following relations. All generators commute, except ∂_i and x_i :

$$[\partial_i, x_i] = \partial_i x_i - x_i \partial_i = 1$$
 for $i = 1, ..., n$.

From PDEs to D-ideals and vice versa

⋄ D gathers linear differential operators with polynomial coefficients

$$P = \sum_{\alpha,\beta \in \mathbb{N}^n} c_{\alpha,\beta} x^{\alpha} \partial^{\beta}, \ c_{\alpha,\beta} \in \mathbb{C} \quad \rightsquigarrow \ \mathsf{PDE} : \boxed{P \bullet f(x_1,\ldots,x_n) = 0}$$

Example: $P = \partial^2 - x \in D$ encodes Airy's equation $f''(x) - x \cdot f(x) = 0$.

left *D*-ideals encode systems of linear PDEs
 operations with *D*-ideals: integral transforms, restrictions, push forward, . . .

Holonomic functions

One variable

A function f(x) is **holonomic** if there exists $P \in D$ that annihilates f, i.e., $P \bullet f = 0$. Multivariate case: $f(x_1, \ldots, x_n)$ is holonomic if $\mathsf{Ann}_D(f)$ is a "holonomic" D-ideal. **Examples:** Feynman integrals, hypergeometric, periods, Airy, some probability distributions, \ldots

Denote by $R_n = \mathbb{C}(x_1, \dots, x_n)\langle \partial_1, \dots, \partial_n \rangle$ the **rational Weyl algebra**.

Theorem (Cauchy–Kovalevskaya–Kashiwara)

Let I be a holonomic D-ideal. The $\mathbb C$ -vector space of holomorphic solutions to I on a simply connected domain in $\mathbb C^n$ outside the singular locus of I has finite dimension

$$|\operatorname{rank}(I)| = \dim_{\mathbb{C}(x_1,...,x_n)} (R_n/R_nI)$$
.

 \Rightarrow A holonomic function is encoded by finite data!

Singularities

D-ideals can be regular singular or irregular singular.

Univariate case: read from growth behavior of general solution near singular points

Example: $\diamond \log(x)$ moderate growth at x = 0 $\diamond \exp(1/x)$ essential singularity at x = 0

Running example

Variables:
$$x_1 = |p_1|^2$$
, $x_2 = |p_2|^2$, $x_3 = |p_1 + p_2|^2$.

The *D*-ideal $I_3(c_0, c_1, c_2, c_3)$

Consider $I_3(c_0, c_1, c_2, c_3) = \langle P_1, P_2, P_3 \rangle \subset D_3$ arising from conformal invariance.

dilatations + conformal boosts

$$p_2$$
 $v_2, k + p_1$
 $v_3, k + p_1 + p_2$
 p_1
 v_1, k

One-loop triangle Feynman diagram with massless propagators and

$$P_1 = 4(x_1\partial_1^2 - x_3\partial_3^2) + 2(2 + c_0 - 2c_1)\partial_1 - 2(2 + c_0 - 2c_3)\partial_3,$$

$$P_2 = 4(x_2\partial_2^2 - x_3\partial_3^2) + 2(2 + c_0 - 2c_2)\partial_2 - 2(2 + c_0 - 2c_3)\partial_3,$$

$$P_3 = (2c_0 - c_1 - c_2 - c_3) + 2(x_3\partial_3 + x_2\partial_2 + x_1\partial_1).$$

Parameters: $c_0 = d$ spacetime dimension

 c_1, c_2, c_3 conformal weights

Choice: $I_3 := I_3(4, 2, 2, 2)$ $\widehat{=}$ conformal ϕ^4 -theory in 4 spacetime dimensions I_3 is regular singular, rank $(I_3) = 4$

Remark: The D-ideal I_3 is the restriction of a GKZ system.

Initial forms

Principal symbol (n = 1)

$$\operatorname{in}_{(0,1)}(x\partial-x^2)=x\xi$$
 is the part of maximal $(0,1)$ -weight $\partial\leadsto\xi$

Several variables: $in_{(0,1)}(x_1\partial_1 + x_2\partial_2 + 1) = x_1\xi_1 + x_2\xi_2$ in general, not a monomial

Algebraically

◆ The **characteristic ideal** of a *D*-ideal *I* is

$$in_{(0,1)}(I) = \langle in_{(0,1)}(P)|P \in I \rangle \subset \mathbb{C}[x_1,\ldots,x_n][\xi_1,\ldots,\xi_n].$$

♦ The characteristic variety of I is

$$\mathsf{Char}(I) \, = \, V(\mathsf{in}_{(0,1)}(I)) \, = \, \big\{ (x,\xi) \, | \, p(x,\xi) = 0 \, \, \mathsf{for \, all} \, \, p \in \mathsf{in}_{(0,1)}(I) \big\} \, \subset \, \mathbb{C}^{2^n} \, .$$

♦ The **singular locus** Sing(*I*) of *I* is the vanishing set of the ideal

$$\left(\,\mathsf{in}_{(0,1)}(\mathit{I}):\langle\xi_1,\ldots,\xi_n\rangle^{(\infty)}\right)\,\cap\,\mathbb{C}[x_1,\ldots,x_n]\,.\quad \mathsf{saturation}\,+\,\mathsf{elimination}$$

Examples

- For $I = \langle x^2 \partial + 1 \rangle \subset D$, $\operatorname{in}_{(0,1)}(I) = \langle x^2 \xi \rangle$ and $\operatorname{Sing}(I) = V(x) = \{0\}$. $\mathbb{C} \cdot \exp(1/x)$
- **9** The characteristic ideal of $I = \langle x_1 \partial_2, x_2 \partial_1 \rangle \subset D_2$ is the $\mathbb{C}[x_1, x_2, \xi_1, \xi_2]$ -ideal $\langle x_1 \xi_2, x_2 \xi_1, x_1 \xi_1 x_2 \xi_2, x_2 \xi_2^2, x_2^2 \xi_2 \rangle$ and $\mathsf{Sing}(I) = V(x_1, x_2) \subset \mathbb{C}^2$. $\mathbb{C} \cdot 1$

Gröbner deformations

Weights of the form
$$(-w, w)$$
, $w = (w_1, ..., w_n) \in \mathbb{R}^n$

- \diamond The **w-weight** of $c_{\alpha,\beta}x^{\alpha}\partial^{\beta}$ is $-w\cdot\alpha+w\cdot\beta$.
- \diamond The **initial form** of $P = \sum c_{\alpha,\beta} x^{\alpha} \partial^{\beta}$ is the subsum of all terms of maximal w-weight.

Initial and indicial ideal (with respect to w)

- ♦ The **initial ideal** of I is the D-ideal $\operatorname{in}_w(I) = \langle \operatorname{in}_{(-w,w)}(P) | P \in I \rangle \subset D$.
- ♦ The **indicial ideal** of I is the $\mathbb{C}[\theta_1, \dots, \theta_n]$ -ideal $\operatorname{ind}_w(I) = R_n \cdot \operatorname{in}_{(-w,w)}(I) \cap \mathbb{C}[\theta_1, \dots, \theta_n]$. $\theta_i =$

 $v_2 = (-1, 2, -1)$ ρ_1 C_1 $V_1 = (1, 0, -1)$ ρ_2 ρ_3

Small Gröbner fan of $I_3 \subset D$, here drawn in $\mathbb{R}^3_w/\mathbb{R}(1,1,1)$.

 $\theta_i = x_i \partial_i$ the *i*-th Euler operator

The zeroes of $\operatorname{ind}_w(I)$ in \mathbb{C}^n are the **exponents** of I. The starting monomials of solutions to I will be of the form $x^A \log(x)^B$ with $A \in V(\operatorname{ind}_w(I))$.

Pipeline: from I to starting terms of series solutions

 $D_{n}\text{-ideal }I \overset{w \in \mathbb{R}^{n}}{\leadsto} \left[\text{in}_{(-w,w)}(I) \right] \rightsquigarrow \left[\text{ind}_{w}(I) \subset \mathbb{C}[\theta_{1},\ldots,\theta_{n}] \right] \overset{V(\text{ind}_{w}(I))}{\leadsto} x^{A} \log(x)^{B}$

Canonical series solutions

Aim: Solutions to
$$I$$
 of the form $F_k(x) = x^A \cdot \sum_{\substack{0 \le p \cdot w \le k, p \in C_{\mathbb{Z}}^* \\ 0 \le b_j < \mathsf{rank}(I)}} c_{pb} x^p \log(x)^b$.

Initial series

The **w-weight** of a monomial $x^A \log(x)^B$ is the real part of $w \cdot A$. The **initial series** $\operatorname{in}_w(f)$ of a function $f = \sum_{A,B} c_{AB} x^A \log(x)^B$ is the subsum of all terms of minimal w-weight.

Proposition

If I is regular holonomic and w a generic weight for I, there exist $\operatorname{rank}(I)$ many canonical series solutions of I which lie in the **Nilsson ring** $N_w(I)$ of I with respect to w,

$$N_w(I) := \mathbb{C}[\![C_w(I)_{\mathbb{Z}}^*]\!][x^{e^1},\ldots,x^{e^r},\log(x_1),\ldots,\log(x_n)].$$

- $\diamond \ \ C_w(I)^*$ the dual cone of the Gröbner cone of $w \qquad \diamond \ \ C_w(I)^*_{\mathbb{Z}} = \ C_w(I)^* \cap \mathbb{Z}^n$
- $\diamond \{e^1, \dots, e^r\}$ the exponents of I

Monomial ordering \prec_w refining w-weight: The number of solutions to I with starting monomial of the form $x^A \log(x)^B$ is the multiplicity of A as zero of $\operatorname{ind}_w(I)$.

M. Saito, B. Sturmfels, and N. Takayama. Gröbner Deformations of Hypergeometric Differential Equations, volume 6 of Algorithms and Computation in Mathematics. Springer, 2000.

The SST algorithm

Theorem (Saito-Sturmfels-Takayama)

Let I be a regular holonomic $\mathbb{Q}[x_1,\ldots,x_n]\langle\partial_1,\ldots,\partial_n\rangle$ -ideal and $w\in\mathbb{R}^n$ generic for I. Let I be given by a Gröbner basis for w. There exists an algorithm which computes all terms up to specified w-weight in the canonical series solutions to I with respect to \prec_w .

Procedure

Input: A regular holonomic D_n -ideal I, its small Gröbner fan Σ in \mathbb{R}^n , a weight vector $w \in \mathbb{R}^n$ that is generic for I, and the desired order $k \in \mathbb{N}$.

... for each starting monomial $x^A \log(x)^B$: solving linear system modulo desired w-weight for vector spaces of monomials of same w-weight. recurrence relations

Output: The canonical series solutions of I with respect to w, truncated at w-weight k.

Starting monomials for I_3

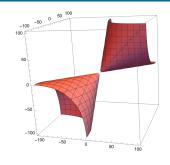
The **singular locus** of I_3 is

$$\mathsf{Sing}\left(I_{3}\right) = V\left(x_{1}x_{2}x_{3} \cdot \lambda\right) \subset \mathbb{C}^{3}.$$

Vanishing locus of the Källén polynomial

$$\lambda = x_1^2 + x_2^2 + x_3^2 - 2(x_1x_2 + x_1x_3 + x_2x_3)$$

 $+ \ coordinate \ hyperplanes \ \{x_i = 0\}$



Initial and indicial ideal for $w=(-1,0,1)\in \mathcal{C}_1$

$$\diamond \ \, \mathsf{in}_{(-w,w)}(\mathit{I}_{3}) \, = \, \, \langle x_{1}\partial_{1} + x_{2}\partial_{2} + x_{3}\partial_{3} + 1, \, x_{2}\partial_{2}^{2} + \partial_{2}, \, x_{3}\partial_{3}^{2} + \partial_{3} \, \rangle \, \subset \, \mathit{D}_{3}$$

$$\diamond \ \operatorname{ind}_{w}(I_{3}) = R_{3} \cdot \operatorname{in}_{(-w,w)}(I) \cap \mathbb{C}[\theta_{1},\theta_{2},\theta_{3}] = \langle \theta_{1} + \theta_{2} + \theta_{3} + 1, \, \theta_{2}^{2}, \, \theta_{3}^{2} \rangle \subset \mathbb{C}[\theta_{1},\theta_{2},\theta_{3}]$$

Exponents of *I*:
$$V(\text{ind}_{w}(I_3)) = \{(-1,0,0)\}.$$
 $= x_1^{-1}x_2^0x_3^0 = 1/x_1$

Change of variables:
$$y_1 = x_1$$
, $y_2 = x_2/x_1$, $y_3 = x_3/x_1$.

Starting monomials of solutions $read\ from\ primary\ decomposition\ of\ ind_w(I)$

$$\diamond 1/y_1 \quad \diamond 1/y_1 \log(y_2) \quad \diamond 1/y_1 \log(y_3) \quad \diamond 1/y_1 \log(y_2) \log(y_3)$$

Canonical series solutions of I_3

Lifting the starting monomials here displayed for f_1 , f_2 , f_3 for w-weight 0 to 4

$$\begin{split} \tilde{f}_1(y_2,y_3) &= 1 + y_2 + y_3 + y_2^2 + 4y_2y_3 + y_3^2 + y_2^3 + 9y_2^2y_3 + y_2^4 + \cdots, \\ \tilde{f}_2(y_2,y_3) &= \log(y_2) + \log(y_2)y_2 + (2 + \log(y_2))y_3 + \log(y_2)y_2^2 + (4 + 4\log(y_2))y_2y_3 \\ &\quad + (3 + \log(y_2))y_3^2 + (\log(y_2))y_2^3 + (6 + 9\log(y_2))y_2^2y_3 + \log(y_2)y_2^4 + \cdots, \\ \tilde{f}_3(y_2,y_3) &= \log(y_3) + (2 + \log(y_3))y_2 + \log(y_3)y_3 + (3 + \log(y_3))y_2^2 \\ &\quad + (4 + 4\log(y_3))y_2y_3 + \log(y_3)y_3^2 + \left(\frac{11}{3} + \log(y_3)\right)y_2^3 \\ &\quad + (15 + 9\log(y_3))y_2^2y_3 + \left(\frac{25}{6} + \log(y_3)\right)y_2^4 + \cdots. \end{split}$$

Then $f_i(x_1,x_2,x_3)=1/x_1\cdot ilde{f_i}(y_2,y_3)$ are canonical series solutions to I_3 . (truncated)

Code

- o in Sage for the bivariate case: https://mathrepo.mis.mpg.de/DModulesFeynman/
- ♦ Macaulay2 package HolonomicSystems M. Sayrafi, C. Berkesch, A. Leykin, H. Tsai

Truncation with respect to w-weight

 $f(x_1, \ldots, x_n)$ general solution of a regular holonomic *D*-ideal *I*

Capturing the weight vector via an auxiliary variable

Choose a generic weight $w \in \mathbb{R}^n$ for I. Set

$$f_w(t,x_1,\ldots,x_n) := f(t^{w_1}x_1,\ldots,t^{w_n}x_n).$$

Merging with canonical series solutions

- **1** From *I*, derive a **Fuchsian system** for $f_w(t, x_1, \dots, x_n)$.
- 2 Solve the system via the path-ordered exponential formalism.
- 3 Compute the asymptotic expansion of $f_w(t,x)$ around t=0:

$$f_w(t,x) = \sum_{k\geq 0} \sum_{m=0}^{m_{\text{max}}} c_{k,m}(x) t^k \log(t)^m.$$

By construction, $c_{k,m}(x)$ has w-weight k.

4 Truncate the expansion at t^k and evaluate at t=1. Nota bene: $f_w|_{t=1} \equiv f$.

T. F. Brown. Iterated Integrals in Quantum Field Theory. In 6th Summer School on Geometric and Topological Methods for Quantum Field Theory, pages 188–240, 2013.

^{2.} W. Wasow. Asymptotic expansions for ordinary differential equations. Pure and Applied Mathematics, Vol. XIV. Interscience Publishers John Wiley & Sons, Inc., New York-London-Sydney, 1965.

Conclusion

In a nutshell

- D-ideals encode crucial properties of their solution functions
 e.g. Feynman integrals, arbitrary loop order, irrespective of whether polylogarithmic, etc.
- algorithmic computation of truncated series solutions by algebraic methods no gauge transform required
- evaluation of solution functions to desired w-weight freedom in choosing weight vector w
- dictionary algebra-physics computing series solutions, Pfaffian system vs. Laporta's algorithm

Thank you for your attention!

J. Henn, E. Pratt, A.-L. S., and S. Zoia. *D*-Module Techniques for Solving Differential Equations in the Context of Feynman Integrals. Preprint arXiv:2303.11105, 2023.

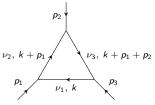
Solutions to I_3

The solution space of I_3 ...

... is spanned by the triangle integral

$$J_{d;\nu_1,\nu_2,\nu_3}^{\rm triangle} = \int_{\mathbb{R}^d} \frac{1}{(-|k|^2)^{\nu_1} \, (-|k+\rho_1|^2)^{\nu_2} \, (-|k+\rho_1+\rho_2|^2)^{\nu_3}} \, \frac{\mathrm{d}^d k}{\mathrm{i} \pi^{\frac{d}{2}}}$$

and its analytic continuations. $rank(I_3) = 4$



One-loop triangle Feynman diagram with massless propagators and massive external particles.

Unitary exponents $\nu_1 = \nu_2 = \nu_3 = 1$, d = 4:

$$\begin{split} f_1(x_1, x_2, x_3) &= J_{4;1,1,1}^{\text{triangle}}(x_1, x_2, x_3) \,, \\ f_2(x_1, x_2, x_3) &= \frac{1}{\sqrt{\lambda}} \log \left(\frac{x_1 - x_2 - x_3 - \sqrt{\lambda}}{x_1 - x_2 - x_3 + \sqrt{\lambda}} \right), \\ f_3(x_1, x_2, x_3) &= \frac{1}{\sqrt{\lambda}} \log \left(\frac{x_2 - x_1 - x_3 - \sqrt{\lambda}}{x_2 - x_1 - x_3 + \sqrt{\lambda}} \right), \\ f_4(x_1, x_2, x_3) &= \frac{1}{\sqrt{\lambda}} \,, \end{split}$$

where $\lambda = x_1^2 + x_2^2 + x_3^2 - 2(x_1x_2 + x_1x_3 + x_2x_3)$ is the **Källén** function.

The conformal group

$$z=\left(z^0,z^1,\ldots,z^{d-1}
ight)^ op$$
 vector of d -dimensional spacetime coordinates $z_1\cdot z_2:=z_1^ op\cdot g\cdot z_2$ $g=\mathrm{diag}(1,-1,\ldots,-1)$ the metric tensor momentum vectors

Translations (Proper) Lorentz transformations Dilatations	$z \longrightarrow z + \epsilon, \ \epsilon \in \mathbb{R}^d$ $z \longrightarrow \Lambda \cdot z, \ \Lambda \in SO(1, d-1)$ $z \longrightarrow \mathrm{e}^\omega z, \ \omega \in \mathbb{R}$
Conformal boosts	$z \longrightarrow \frac{z - z ^2 \epsilon}{1 - 2z \cdot \epsilon + z ^2 \epsilon ^2}, \ \epsilon \in \mathbb{R}^d$

Poincaré group conformal group

symmetry group of Einstein's theory of special relativity Poincaré + dilatations + conformal boosts

Invariance under...

- translations implies momentum conservation
- \diamond Lorentz transformation implies dependency on Mandelstam invariants $p_k \cdot p_\ell$ only

Generators in position space to momentum space via Fourier transform

- $\mathfrak{D}_n = -i \sum_{k=1}^n (z_k \cdot \partial_{z_k} + c_k)$
- \diamond conformal boosts: $\Re_n = i \sum_{k=1}^n \left[|z_k|^2 \partial_{z_k} 2 z_k (z_k \cdot \partial_{z_k}) 2 c_k z_k \right]$

Running example: n = 3, momenta p_1, p_2, p_3 , variables $x_i = |p_i|^2$

- \diamond P_3 stems from $\widehat{\mathfrak{D}_3}$ \diamond
- \diamond P_1, P_2 stem from $\widehat{\mathfrak{K}}_3$

Systems in matrix form

- $\diamond I$ a holonomic D_n -ideal of rank $m = \text{rank}(I), f \in \text{Sol}(I)$
- $\diamond 1, s_2, \dots, s_m$ a $\mathbb{C}(x)$ -basis of R_n/R_nI standard monomials for a Gröbner basis of I

Pfaffian system

Set $F = (f, s_2 \bullet f, \dots, s_m \bullet f)^{\top}$. There exist $P_1, \dots, P_n \in \mathbb{C}(x_1, \dots, x_n)^{m \times m}$ for which

$$\partial_i \bullet F = P_i \cdot F$$
.

The matrices P_i fulfill $P_iP_i - P_iP_i = \partial_i \bullet P_i - \partial_i \bullet P_i$ for all i, j. integrability

If all poles are of order at most 1, the system is Fuchsian. To arrive at a Fuchsian form, one might need a gauge transform. Wasow's method

Construction of a Pfaffian system IBP reduction with Laporta's algorithm

$$\partial^a$$
 Feynman integrals a in ∂^a propagator powers $\partial^a Q_i = 0$ in $R_n/R_n I$ IBP identities $\mathbb{C}(x)$ -basis of $R_n/R_n I$ set of master integrals

^{1.} V. Chestnov, F. Gasparotto, M. K. Mandal, P. Mastrolia, S.-J. Matsubara-Heo, H. J. Munch, N. Takayama. Macaulay matrix for Feynman integrals: Linear relations and intersection numbers. J. High Energy Phys., 187(2022), 2022.

^{2.} W. Wasow. Asymptotic expansions for ordinary differential equations. Pure and Applied Mathematics, Vol. XIV. Interscience Publishers John Wiley & Sons, Inc., New York-London-Sydney, 1965.

The SST algorithm

Input: A regular holonomic D_n -ideal I, its small Gröbner fan Σ in \mathbb{R}^n , a weight vector $w \in \mathbb{R}^n$ that is generic for I, and the desired order $k+1 \in \mathbb{N}$.

- **1** Determine a Gröbner basis $G = \{g_1, \dots, g_d\}$ of I with respect to w.
- ② Write each $g \in G$ as $x^{\alpha}g = f h$ with $\alpha \in \mathbb{Z}^n$ such that $f \in \mathbb{K}[\theta_1, \dots, \theta_n]$ and $h \in \mathbb{K}[x_1^{\pm 1}, \dots, x_n^{\pm 1}]\langle \partial_1, \dots, \partial_n \rangle$ with $\operatorname{ord}_{(-w,w)}(h) < 0$.
- **3** Compute the indicial ideal ind_w(I) and its rank(I) many solutions. They are the form $x^A \log(x)^B$ with $A \in V(\operatorname{ind_w}(I))$. For each starting of these monomials, carry out Step 4.
- Assume the partial solution

$$F_k(x) = x^A \cdot \sum_{0 \le p \cdot w \le k, \ p \in C_{\mathbb{Z}}^*} c_{pb} x^p \log(x)^b.$$

is known. Solve the linear system

$$\begin{split} &(f_1,\ldots,f_d) \bullet E_{k+1}(x) = (h_1-f_1,\ldots,h_d-f_d) \bullet F_k(x) \text{ mod } w\text{-weight } k+2 \\ \text{for } E_{k+1} \in \sum_{p \cdot w = k+1, \ p \in C_{\mathbb{Z}}^*} \mathcal{L}'_p \text{ of } w\text{-weight } k+1. \quad \text{Adding } E_{k+1} \text{ to } F_k \text{ lifts } F_k \text{ to } F_{k+1}. \\ \mathcal{L}'_p \text{ the subspace of } \mathcal{L}_p = x^A \sum_{0 \leq b_i \leq \text{rank}(I)} \mathbb{K} \cdot x^p \log(x)^b \text{ spanned by monomials } \notin \text{Start}_{\prec_w}(I) \end{split}$$

Output: The canonical series solutions of I with respect to w, truncated at w-weight k+1.

SST algorithm: a hypergeometric example

Consider the *D*-ideal *I* generated by $P = \theta(\theta - 3) - x(\theta + a)(\theta + b)$.

- 1 I is holonomic of rank $ord_{(0,1)}(P) = 2$.
- **2** Gröbner fan of I: two maximal cones $\pm \mathbb{R}_{\geq 0}$.
- **3** For the weight w = 1, $\operatorname{in}_{(-w,w)}(I) = \langle \theta(\theta 3) \rangle = \operatorname{ind}_w(I)$.
- **4** Exponents of I: $V(\operatorname{ind}_w(I)) = \{0,3\}$. starting monomials $x^0 = 1$ and x^3
- **6** Choose $\frac{x^3}{x^3}$ as starting monomial, $L_p = \mathbb{C} \cdot \{x^{p+3}, x^{p+3} \log(x)\}$. $\left\{x^3 \sum_p c_{p,1} x^p + c_{p,2} x^p \log(x)\right\}$
- **6** Write P = f h, where $f = \theta(\theta 3)$ and $h = x(\theta + a)(\theta + b)$. Action of θ on L_p :

$$\theta \bullet x^{p+3} \, = \, (p+3) x^{p+3} \quad \text{and} \quad \, \theta \bullet (x^{p+3} \log(x)) \, = \, x^{p+3} + (p+3) x^{p+3} \log(x) \, .$$

Thus, the matrix of the operator θ in the basis $\{x^{p+3}, x^{p+3} \log(x)\}$ is

$$\begin{bmatrix} p+3 & 1 \\ 0 & p+3 \end{bmatrix}.$$

7 Let $c_{p,1}$ and $c_{p,2}$ be the coefficients of x^{p+3} and $x^{p+3}\log(x)$ in the power series expansion. Then we can write our operators as matrices, and our **recurrence** as

$$\begin{bmatrix} p & 1 \\ 0 & p \end{bmatrix} \begin{bmatrix} p+3 & 1 \\ 0 & p+3 \end{bmatrix} \begin{bmatrix} c_{p,1} \\ c_{p,2} \end{bmatrix} = \begin{bmatrix} p-a+2 & 1 \\ 0 & p-a+2 \end{bmatrix} \begin{bmatrix} p-b+2 & 1 \\ 0 & p-b+2 \end{bmatrix} \begin{bmatrix} c_{p-1,1} \\ c_{p-1,2} \end{bmatrix}$$

with initial values $c_{0,1} = 1$, $c_{0,2} = 0$. Solving the recurrence yields

$$c_{p,1}=0$$
 and $c_{p,2}=rac{(a+3)_p(b+3)_p}{(1)_p(4)_p}$.